

Taylor Methods and the Dynamical Analysis of Muon Accelerators

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What is a Particle Accelerator?

- Guide charged particles along a particular path, using magnets
- That path is often circular, and the particles make many (10^8 !) turns
- The particles are often accelerated from lower energies to higher energies
- Usually trying to do something with these particles eventually
 - ◆ Generally collide them into each other
 - ◆ Sometimes collide them into a stationary object
 - ◆ Sometimes make use of the radiation that they produce
 - ◆ For muons, sometimes we leave them in a ring and let them decay, producing a beam of neutrinos

- Define reference curve $\mathbf{x}_0(s)$, s is arc length
- Coordinates are defined as deviations from this reference curve

$$\hat{\mathbf{s}}(s) = \frac{d\mathbf{x}_0}{ds} \quad \hat{\mathbf{x}}(s) \cdot \hat{\mathbf{s}}(s) = 0 \quad \hat{\mathbf{y}}(s) = \hat{\mathbf{s}}(s) \times \hat{\mathbf{x}}(s)$$

$$x = [\mathbf{z} - \mathbf{x}_0(s)] \cdot \hat{\mathbf{x}} \quad y = [\mathbf{z} - \mathbf{x}_0(s)] \cdot \hat{\mathbf{y}}$$

- Hamiltonian dynamics
 - ◆ Independent variable is s
 - ◆ Hamiltonian, function of phase space coordinates (x, p_x, y, p_y, t, E)
 - ◆ Vector field based on derivatives of Hamiltonian:
 $(\partial_{p_x} H, -\partial_x H, \partial_{p_y} H, -\partial_y H, -\partial_E H, \partial_t H)$

$$- (1 + hx) A_s$$

$$- (1 + hx) \sqrt{\left(\frac{E - q\phi}{c}\right)^2 - (p_x - qA_x)^2 - (p_y - qA_y)^2 - (mc)^2}$$

- Simplified: ignoring torsion and vertical bending
- h is curvature of reference curve
- Electrostatic potential ϕ , magnetic vector potential (A_x, A_y, A_s)
 - ◆ Generally $\phi = 0$

- Vector potentials obey Maxwell's equations; for no curvature, time-independent magnetic fields:

$$A_s(x, y, s) =$$

$$\sum_{m=1} \sum_{k=0} \frac{1}{k!(k+m)!} \Re\{A_m^{(2k)}(s)(x+iy)^m\} \left(-\frac{x^2+y^2}{4}\right)^k$$

$$A_x(x, y, s) =$$

$$\sum_{m=1} \frac{1}{2m} \sum_{k=0} \frac{1}{k!(k+m)!} \Re\{A_m^{(2k+1)}(s)(x+iy)^m\} \left(-\frac{x^2+y^2}{4}\right)^k$$

$$A_y(x, y, s) =$$

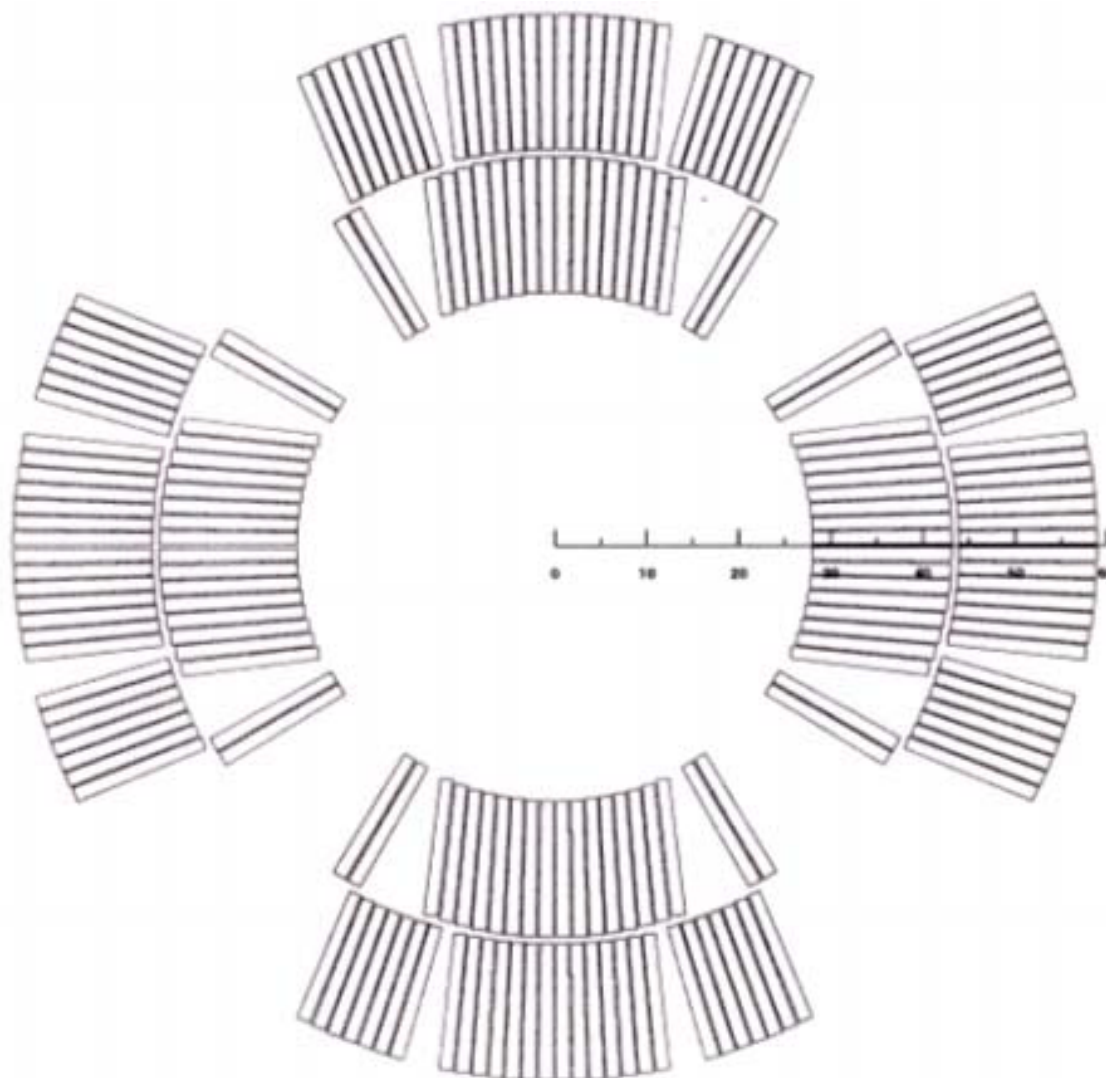
$$\sum_{m=1} \frac{1}{2m} \sum_{k=0} \frac{1}{k!(k+m)!} \Im\{A_m^{(2k+1)}(s)(x+iy)^m\} \left(-\frac{x^2+y^2}{4}\right)^k$$

- For given m , lowest order contribution to vector potential is

$$A_s(x, y, s) = \Re\{A_m(s)(x + iy)^m\}$$

- When bending particle, use $A_s = -A_1(s)x$, with A_1 and h chosen together to keep particle on (near) reference curve (dipole)
 - ♦ I'm cheating, ignoring curvature terms here
- To give linear stability of orbits about reference curve, use term $A_s = A_2(s)(x^2 - y^2)/2$ (quadrupole)
 - ♦ Control other things we this also, discussed shortly
- Notice: these are low-order Taylor series, they generate low-order Taylor series in the Hamiltonian and vector field
 - ♦ If there were no s dependence, there would be no higher-order terms
- These A_s create linear terms in the vector field

Quadrupole Cross-Section



- Field in magnet is designed to be as constant longitudinally as possible



- Short sections on ends where field is not constant, giving higher-order terms
- The integral of A_m gives the lowest-order effect
- The longer the magnet, the smaller the relative additional contribution of the ends is
 - ◆ End length is proportional to magnet aperture
 - ◆ All these effects generate higher-order terms
- A_x and A_y are higher-order

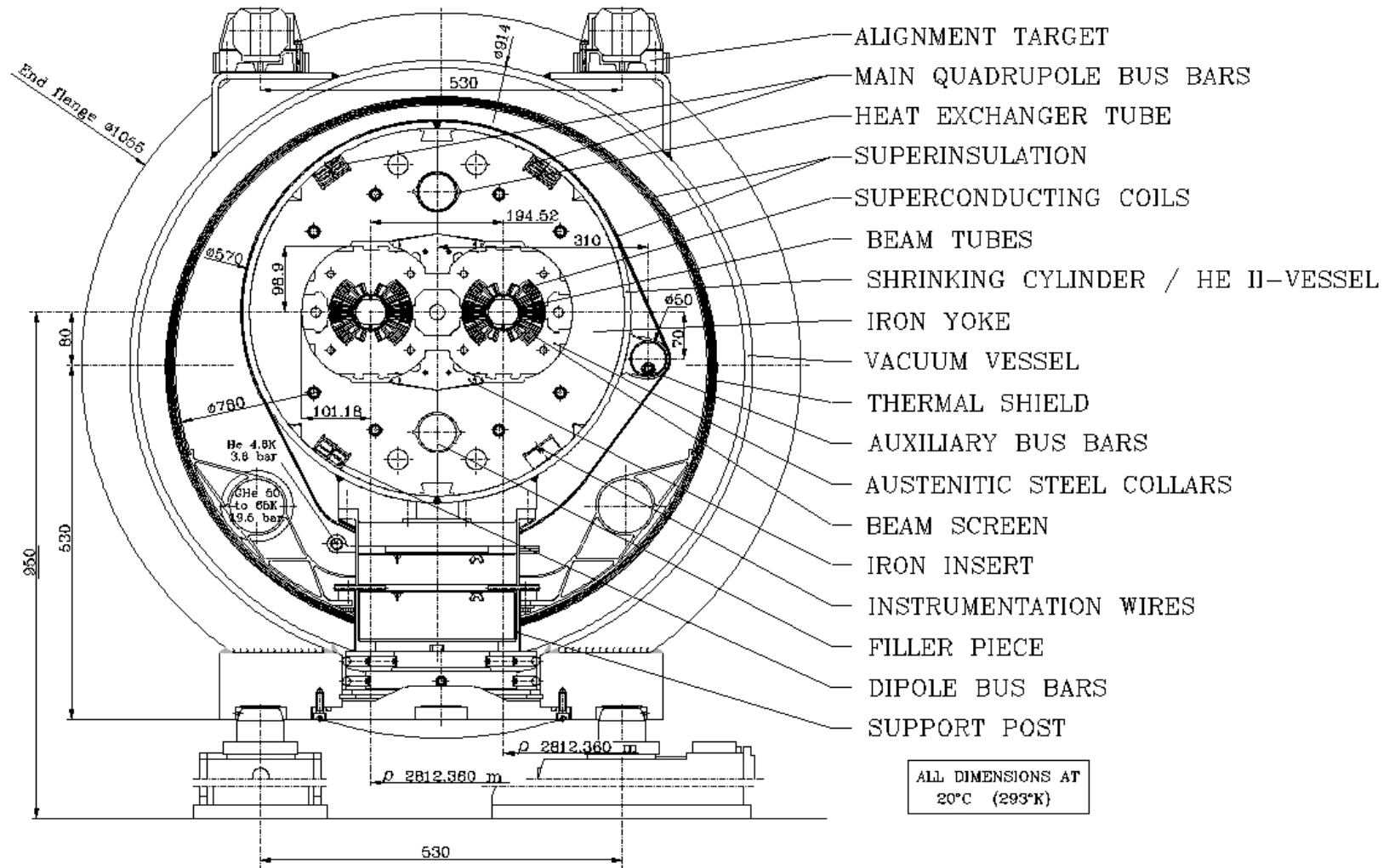
- Nominal dipole length is 14.312 m
- Inner coil diameter: 5.6 cm; outer coil diameter: about 11.8 cm.
- Thus, end length is small fraction of total length, leading to small relative importance of ends
- There are specific magnets for which ends are important, but they are small in number
- The RMS beam size in the arc is 1.7 mm
- There are higher multipole components in these magnets (A_m , with $m > 2$), which are kept to about 10^{-4} relative to the desired multipole component at a radius of 1.7 cm.
- Hamiltonian is linear in A_s

- There are nonlinear magnets (A_m with $m > 2$) intentionally used, but they should give small contributions
- Conclusion:
 - ◆ Contributions of magnets to the vector field are primarily linear
 - ◆ Higher-order contributions are small
 - ◆ Magnetic contributions to map are well-represented by a rapidly converging Taylor series

LHC Magnet Photo



LHC Magnet Cross-Section



- Can expand square root part of Hamiltonian (ignore potentials, curvature) ($p = \sqrt{(E/c)^2 - (mc)^2}$)

$$-p + \frac{p_x^2 + p_y^2}{2p} + \frac{(p_x^2 + p_y^2)^2}{8p^3} + \dots$$

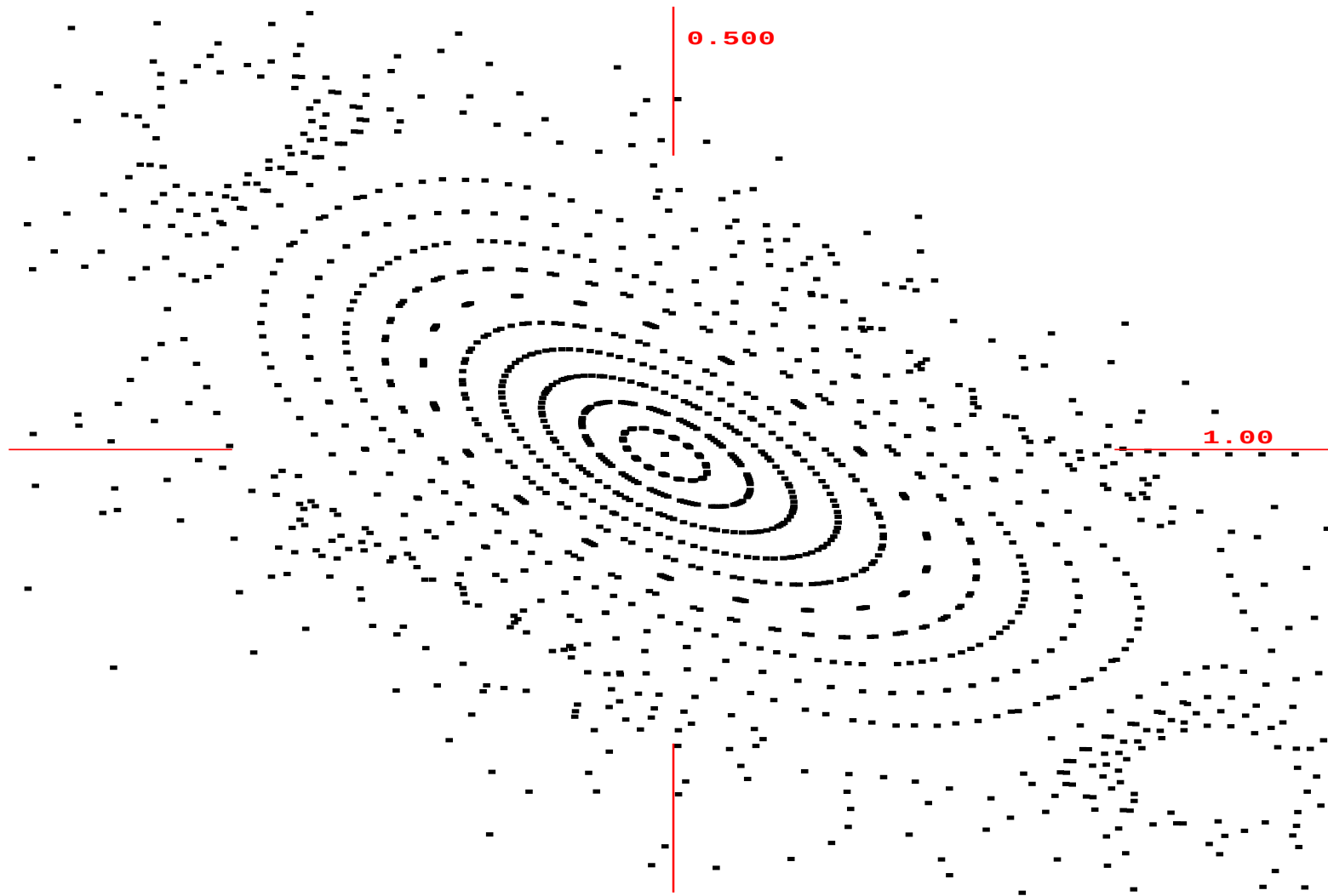
- Second term gives linear contribution to vector field
- Subsequent terms gives terms that are higher order in p_x/p and p_y/p
- LHC: the largest value for p_x/p is 1.2×10^{-4}
- Conclusion: square root contribution is well-represented by a rapidly converging Taylor series

- Say we have a reference energy E_0 ; the energy is $E_0 + \Delta$
- Energy appears in above square root term as $\sqrt{(E/c)^2 - (mc)^2}$
- LHC: RMS $\Delta/pc = 3.06 \times 10^{-4}$
- Again, expect a rapidly converging Taylor series

- Terms like $1 + hx$ appear in Hamiltonian, h is curvature
- LHC: $h = 3.57 \times 10^{-4} \text{ m}^{-1}$, typical $x = 1.7 \text{ mm}$
- Similar curvature terms appear in magnetic field expansions
- Again, terms are small, Taylor series converges rapidly

- For most high energy particle accelerators
 - ◆ Vector field is well represented *globally* by a Taylor series
 - ◆ Terms in the Taylor series converge rapidly due to small amplitudes
 - ◆ Magnetic fields look like a Taylor series, by design
 - ◆ Linearized vector field gives a very good picture of the dynamics
- Cannot ignore nonlinear terms
 - ◆ With only linear terms, motion is stable for arbitrary displacements from the reference orbit
 - ◆ Due to small amplitudes, typically don't need an extremely high order Taylor series to represent the motion
 - ◆ Typically 8th order is sufficient for large machines
- Exceptions to this
 - ◆ Small rings: angles large, curvature large
 - ◆ Muon machines

Dynamic Aperture Plot



- Replace integration of equations of motion with evaluation of Taylor series
 - ◆ LHC contains around 2500 magnets
 - ◆ Beam makes around 6×10^8 turns
 - ◆ Integrating through each magnet can be prohibitive
 - ◆ Instead, construct Taylor series representing one-turn map
- Design an analysis: linear
 - ◆ Stability of fixed point
 - ◆ Beam sizes
- Design and analysis: nonlinear
 - ◆ Amplitude-dependent tunes (eigenvalues)
 - ◆ Resonances (higher-order fixed points away from reference orbit)

- Beam sizes, energy spreads are large
- Different types of lattices
 - ◆ Cooling lattices
 - ◆ Fixed Field Alternating Gradient (FFAG) acceleration

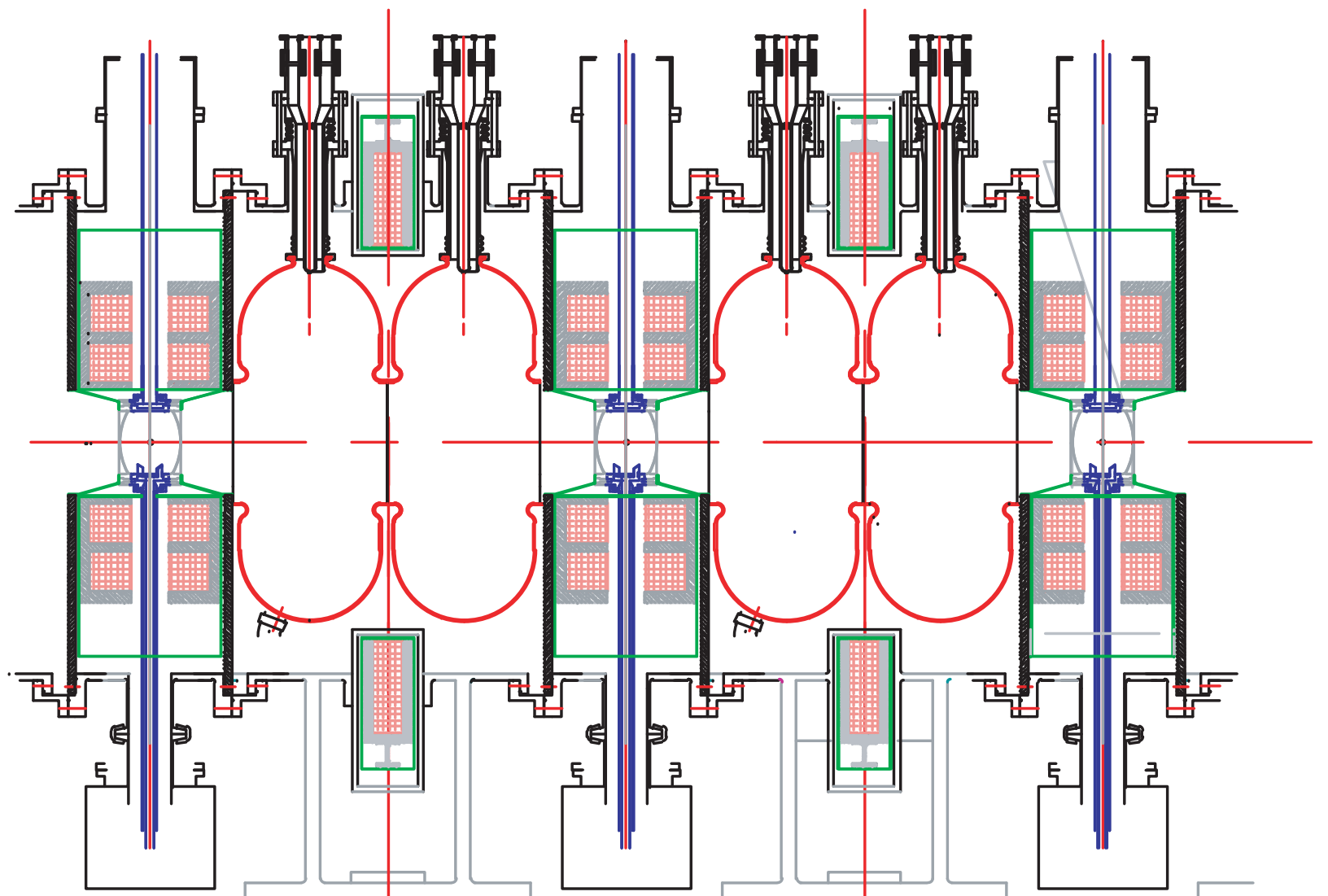
- Cooling cells

- ◆ Cooling requires extremely large angular spreads: as high as $p_x/p = 0.3$
- ◆ Lattices may transport a factor of 2 in energy
- ◆ Lattices must be compact to achieve these acceptances
- ◆ Large beam size requires large magnet aperture
 - ★ Beam fills significant fraction of aperture
- ◆ Field no longer mostly constant
 - ★ Higher order terms are significant
 - ★ Series does not converge as quickly

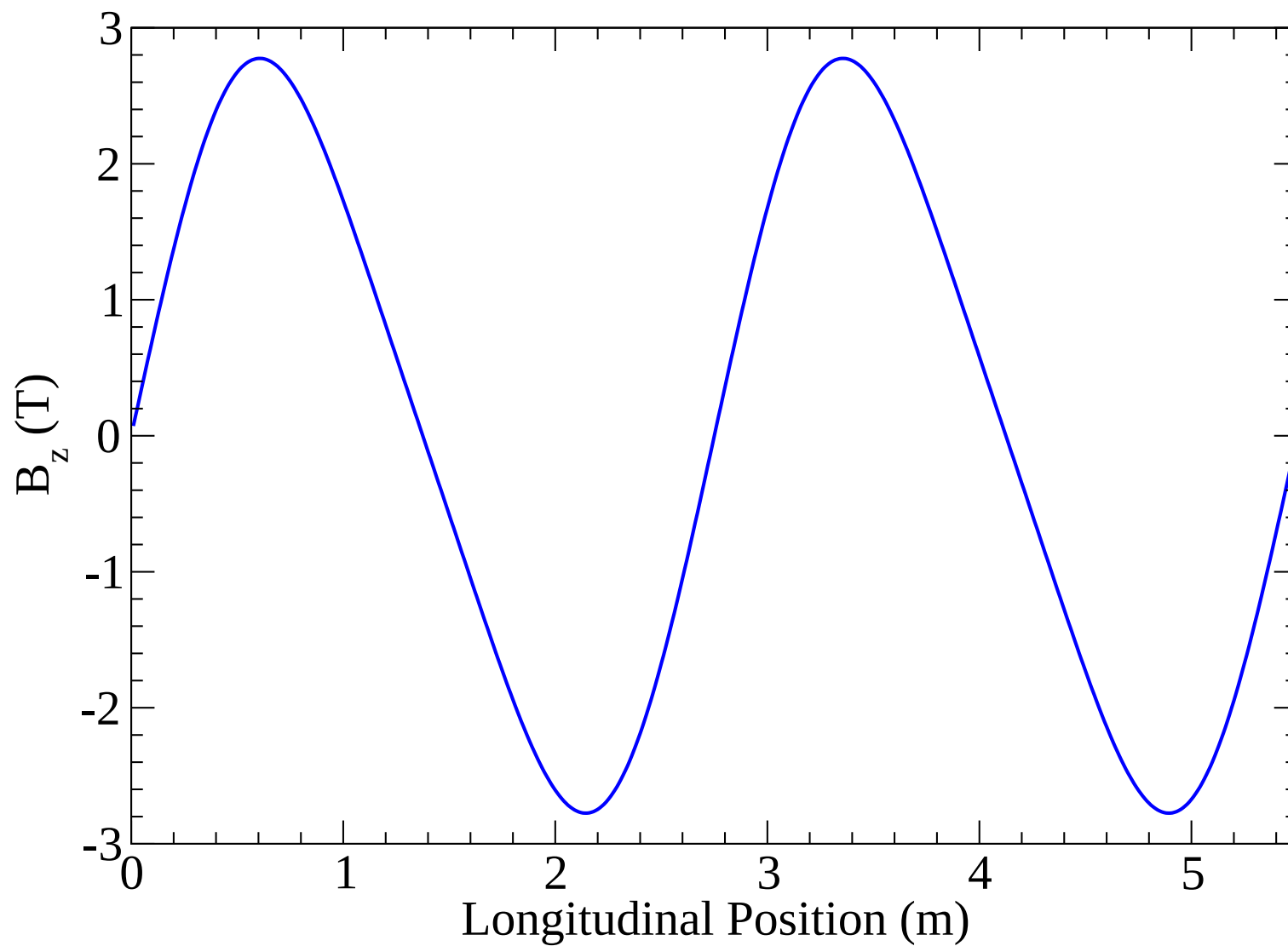
- FFAG lattices

- ◆ Normally when accelerating in a ring, increase magnetic fields with momentum: equations of motion (nearly) invariant
- ◆ FFAG: keep fields fixed. Arc must accept a factor of 2 or more in energy

Cooling Cell

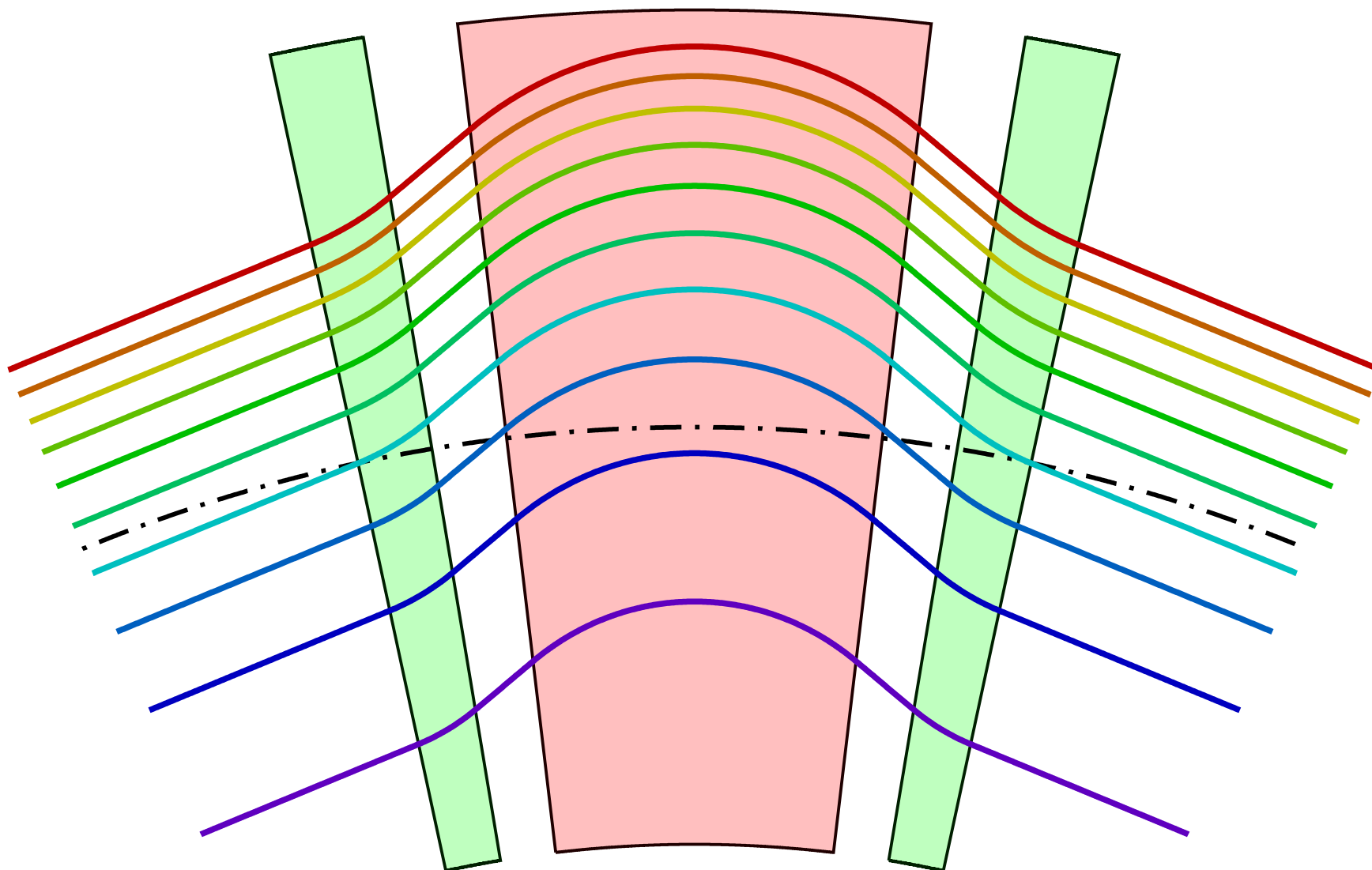


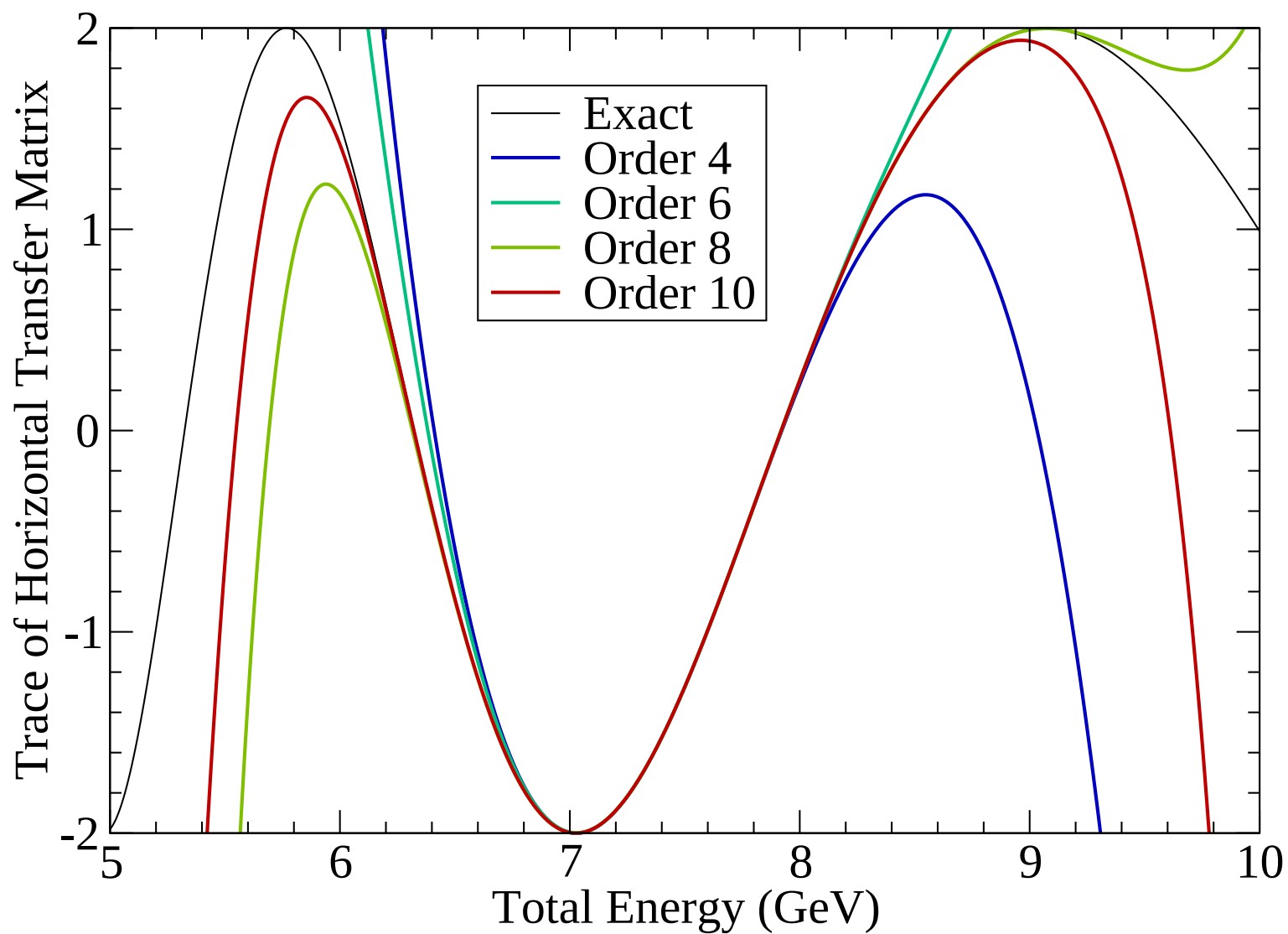
Field in Cooling Cell



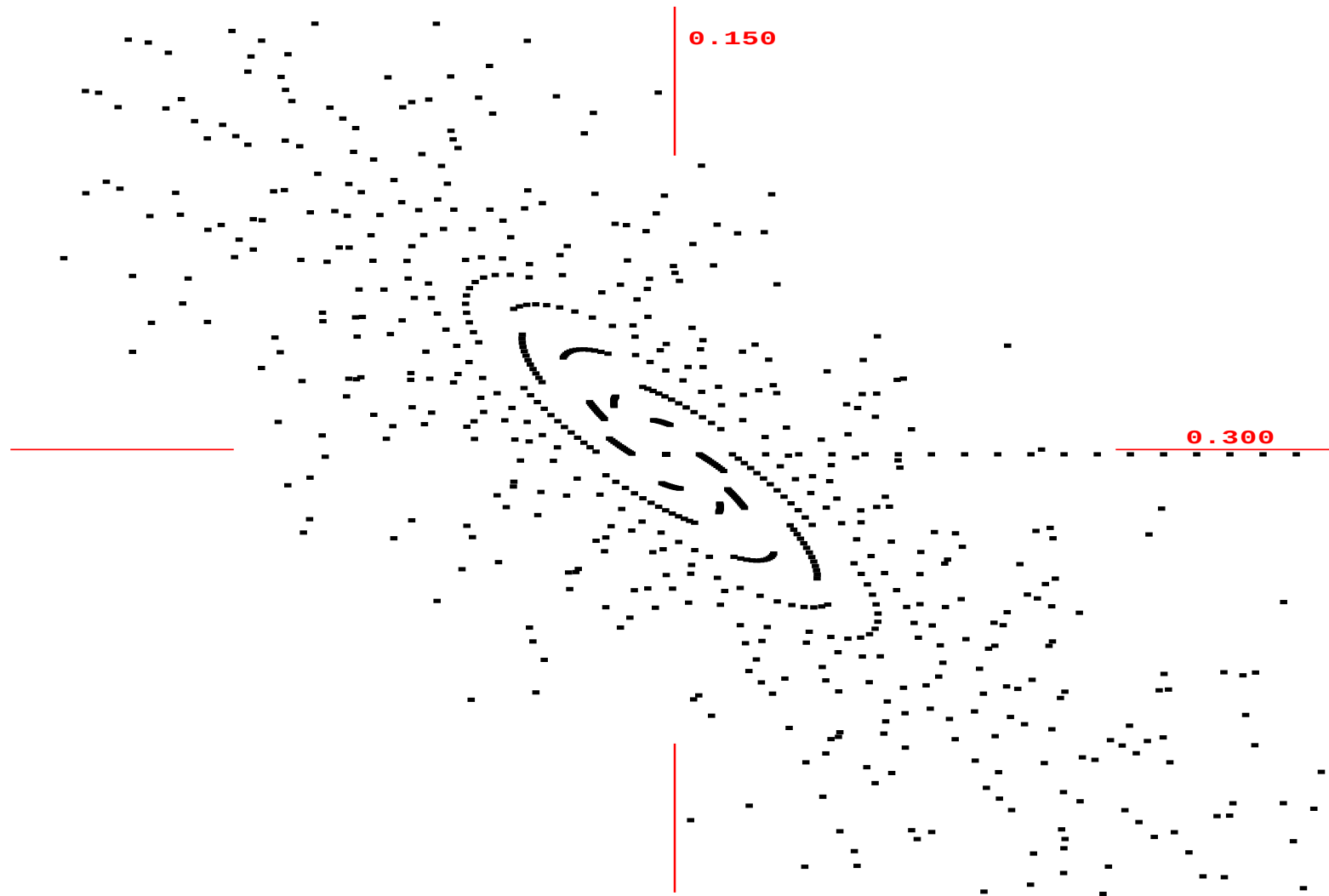
- Analysis of lattice
 - ◆ Find fixed point of map for fixed Δ
 - ◆ The closed orbit is the phase space variables as a function of s passing through this fixed point
 - ◆ Translate coordinates to coordinates relative to this closed orbit
 - ◆ Write map in these new coordinates
 - ◆ Now have eigenvalues as a function of Δ , other linear lattice parameters
- Problem: truncation. Evaluating Δ at large values gives effectively lower map in transverse variables
 - ◆ Linear lattice functions incorrect for large Δ
 - ◆ Tracking results manifestly wrong
- Solutions do not always converge with increasing order

Closed Orbits





Off-Energy Tracking



- Need separate parameterization in Δ
 - ◆ Some kind of spline function
 - ◆ Don't make total order count order of polynomial in Δ
 - ◆ Domain decomposition
- Potentially have similar problems for large angles, highly nonlinear magnets, etc.

- Traditional accelerators seem to be designed to be represented by Taylor series
 - ◆ Idealized contribution of a magnet is a Taylor series
 - ◆ Most deviations are small, making the Taylor series converge quickly
 - ◆ A single Taylor series gives a global map representation
- In muon accelerators, deviations are not small
 - ◆ Leads to problems representing large energy ranges
 - ★ To use Taylor methods, must parametrize energy differently
 - ◆ May also encounter problems resulting from
 - ★ Large angles
 - ★ Magnets whose end fields are significant