

# The Multiple Level Fast Multipole Method in the Differential Algebra Framework for Space Charge Field Calculation and Simulations on femtosecond electron imaging and spectroscopy

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- Multiple level fast multipole algorithm (MLFMA) in the differential algebra framework
- Simulation results on femtosecond electron imaging and spectroscopy

## I. Multiple Level Fast Multipole Algorithm in DA framework

- Space charge effect is an important effect in the photo-emission process that generates the femtosecond electron bunch
- To simulate space charge effect, one needs an algorithm that has **good efficiency** and **good accuracy**.
- The MLFMA is a good choice, because it
  - scales **linearly** with the number of particles for any **arbitrary charge distribution (grid free)**
  - does NOT artificially smooth the field
- Using DA, one can
  - calculate not only the field, but also its high order derivatives
  - represent the multipole/local expansions in Cartesian coordinates naturally
  - considerably simplify the math

- The DA based MLFMA is developed as a package for COSY Infinity
- COSY is a scientific computing system
  - DA and TM data types (ODE, flows, PDE, automatic differentiation, range bounding, etc.)
  - beam physics package
- More information on <http://www.cosyinfinity.org>

Two operations in COSY:

- Automatic Taylor expansion of a function

$$f(x + \delta x) = f(x) + f'(x)\delta x + \frac{1}{2!}f''(x)\delta x^2 + \frac{1}{3!}f'''(x)\delta x^3 + \dots$$

In COSY,

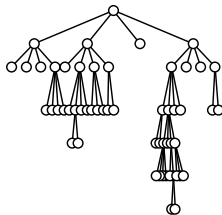
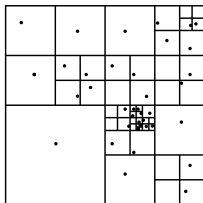
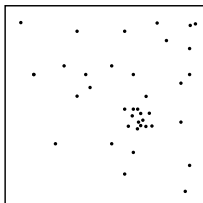
$$f(x + da(1)) = f(x) + f'(x)da(1) + \frac{1}{2!}f''(x)da(1)^2 + \frac{1}{3!}f'''(x)da(1)^3 + \dots$$

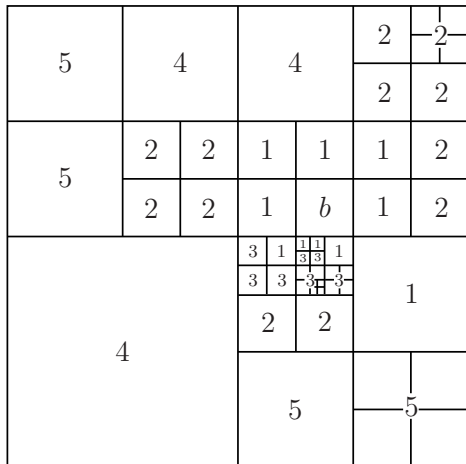
- Composition of two maps

$$G(x) = G(F) \circ F(x), \text{ or } G(x) = G(F(x))$$

In COSY, it can be done by the command POLVAL.

- **Fast Multipole Method (FMM)**, L.Greengard and V.Rokhlin, 1987
- Cut the whole region into boxes of **hierarchical tree structure**.
- **Multipole expansions** and **local expansions**.
- For any arbitrary distribution, **scales linearly** with the number of particles.

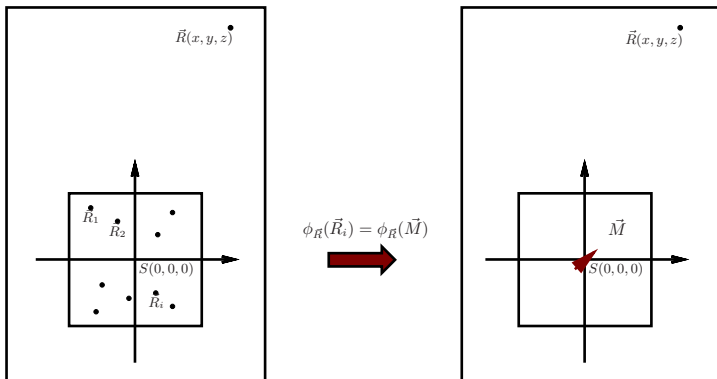




Relations		Operations
$c \in 1_b$	$b \in 1_c$	$C_c \rightarrow C_b, C_b \rightarrow C_c$
$c \in 2_b$	$b \in 2_c$	$M_c \rightarrow L_b, M_b \rightarrow L_c$
$c \in 3_b$	$b \in 4_c$	$M_c \rightarrow C_b, C_b \rightarrow L_c$
$c \in 4_b$	$b \in 3_c$	$C_c \rightarrow L_b, M_b \rightarrow C_c$
$c \in 5_b$	$b \in 5_c$	Do nothing



## Multipole expansion from charges (for the childless boxes)



$$\begin{aligned}
 \phi &= \sum_{i=1}^n \frac{q_i}{\sqrt{(x-x_i)^2 + (y-y_i)^2 + (z-z_i)^2}} \\
 &= \sum_{i=1}^n \frac{d_r \cdot q_i}{\sqrt{1 + (x_i^2 + y_i^2 + z_i^2)d_r^2 - 2x_id_x - 2y_id_y - 2z_id_z}} \\
 &= d_r \cdot \bar{\phi}_{c2m}
 \end{aligned}$$

with

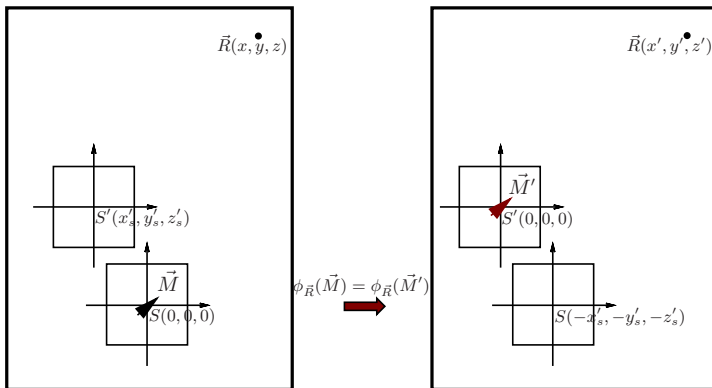
$$d_x = \frac{x}{x^2 + y^2 + z^2}, \quad d_y = \frac{y}{x^2 + y^2 + z^2},$$

$$d_z = \frac{z}{x^2 + y^2 + z^2}, \quad d_r = \sqrt{d_x^2 + d_y^2 + d_z^2},$$

$$\bar{\phi}_{c2m} = \sum_{i=1}^n \left\{ q_i / \sqrt{1 + (x_i^2 + y_i^2 + z_i^2)d_r^2 - 2x_id_x - 2y_id_y - 2z_id_z} \right\}.$$

Error  $|\epsilon| \leq C \cdot \left(\frac{a}{r}\right)^{p+1} \cdot \frac{1}{r-a}$ , where  $C = \sum_{i=1}^n |q_i|$  and  $r_i \leq a$  for any  $i$ .

Translate the position of a multipole expansion



In parent box frame, new DA variables are chosen as

$$d'_x = \frac{x - x'_o}{r'^2} = \frac{x'}{r'^2}, \quad d'_y = \frac{y - y'_o}{r'^2} = \frac{y'}{r'^2}$$

$$d'_z = \frac{z - z'_o}{r'^2} = \frac{z'}{r'^2},$$

Relation between the old and new DA variables.

$$\begin{aligned} d_x &= (d'_x + x'_o \cdot (d'^2_x + d'^2_y + d'^2_z)) \cdot R, \\ d_y &= (d'_y + y'_o \cdot (d'^2_x + d'^2_y + d'^2_z)) \cdot R, \\ d_z &= (d'_z + z'_o \cdot (d'^2_x + d'^2_y + d'^2_z)) \cdot R, \end{aligned} \quad Mm2m$$

with

$$R = \frac{1}{1 + (x'^2_o + y'^2_o + z'^2_o)(d'^2_x + d'^2_y + d'^2_z) + 2x'_o d'_x + 2y'_o d'_y + 2z'_o d'_z}.$$

In child box frame  $\phi = d_r \cdot \bar{\phi}_{c2m}$ .

In the parent box frame

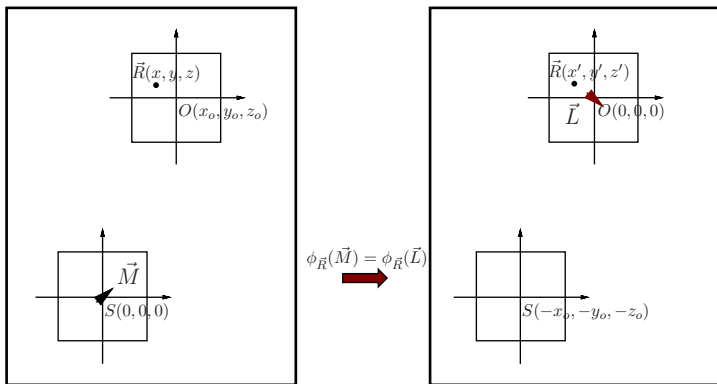
$$\phi' = d'_r \cdot \sqrt{R} \cdot \phi_{m2m} = d'_r \cdot \bar{\phi}_{m2m}$$

with  $d'_r = \sqrt{d_x'^2 + d_y'^2 + d_z'^2}$ ,

and  $\phi_{m2m} = \bar{\phi}_{c2m} \circ M_{m2m}$ ,

where  $M_{m2m}$  is the map from the old DA variables into the new DA variables

Convert a multipole expansion into a local expansion



New DA variables in the observer frame

$$d'_x = x - x'_o = x',$$

$$d'_y = y - y'_o = y',$$

$$d'_z = z - z'_o = z'.$$

The relation between the new and the old DA variables

$$\begin{aligned} d_x &= (x'_o + d'_x) \cdot R, \\ d_y &= (y'_o + d'_y) \cdot R, \\ d_z &= (z'_o + d'_z) \cdot R. \end{aligned} \quad M_{m2l}$$

with

$$R = \frac{1}{(x'_o + d'_x)^2 + (y'_o + d'_y)^2 + (z'_o + d'_z)^2}.$$

The multipole expansion in the source frame  $\phi = d_r \cdot \bar{\phi}$   
 The local expansion in the observer frame

$$\phi = \sqrt{R} \cdot \bar{\phi}_{m2l} = \phi_{m2l}$$

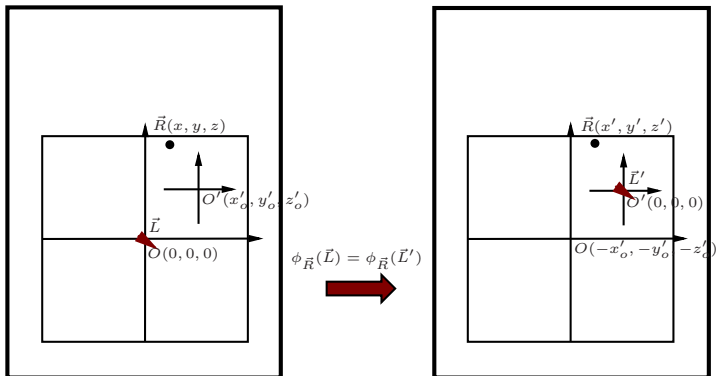
where  $\sqrt{R}$  is converted from  $d_r, \bar{\phi}_{m2l} = \bar{\phi} \circ M_{m2l}$ ,  
 and  $M_{m2l}$  is the map between the DA variables.

Error

$$|\epsilon| \leq C \cdot \left(\frac{a}{r'_o}\right)^{p+1} \cdot \frac{1}{r'_o - a} + C \cdot \left(\frac{r'}{b}\right)^{p+1} \cdot \frac{1}{b - r'}$$




Translate a local expansion from a parent box to its child boxes



DA variables in the child box frame

$$\begin{aligned} d_x &= x'_o + d'_x, \\ d_y &= y'_o + d'_y, \\ d_z &= z'_o + d'_z. \end{aligned}$$

$M_{l2l}$

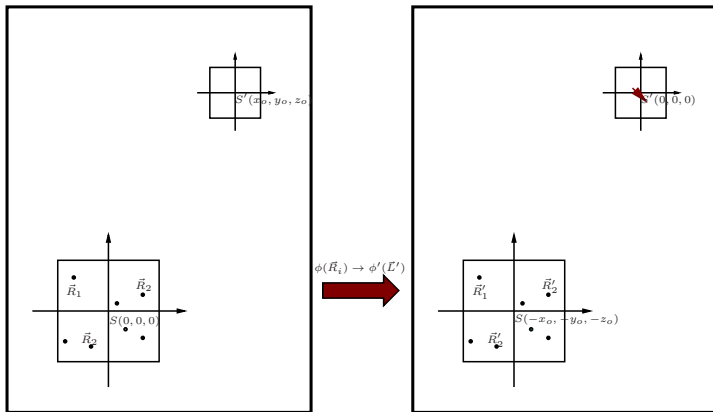


The local expansion in the parent box frame is  $\phi_{m2l}$ .  
The local expansion in the child box frame is

$$\phi = \phi_{m2l} \circ M_{l2l} = \phi_{l2l},$$

where  $M_{l2l}$  is the map between the old and the new DA variables.

Calculate the local expansion from charges.



In the observer (small box) frame, the new DA variables are

$$\begin{aligned}d'_x &= x - x'_o = x', \\d'_y &= y - y'_o = y', \\d'_z &= z - z'_o = z'.\end{aligned}$$

Then the local expansion is

$$\begin{aligned}\phi_L &= \sum_{i=1}^n \frac{q_i}{\sqrt{(x - x_i)^2 + (y - y_i)^2 + (z - z_i)^2}} \\ &= \sum_{i=1}^n \frac{q_i}{\sqrt{(x'_o - x_i + d'_x)^2 + (y'_o - y_i + d'_y)^2 + (z'_o - z_i + d'_z)^2}}\end{aligned}$$

Error

$$|\epsilon| \leq C \cdot \left(\frac{r'}{b}\right)^{p+1} \cdot \frac{1}{b - r'}$$

- Now we have the potential expressed as a **polynomial of coordinates** up to order  $p$ .
- Take the derivative of a coordinates to get the field expression in a **polynomial of coordinates** up to order  $p - 1$ .
- Submit the charge positions into the expression to calculate the potential/field.

Calculate the field from the multipole expansion.

The multipole expansion is  $\phi = d_r \cdot \bar{\phi}$ , then

$$E_x = \left\{ -\frac{\partial \bar{\phi}}{\partial d_x} \cdot (d_r^2 - 2d_x^2) + 2\frac{\partial \bar{\phi}}{\partial d_y} \cdot d_x d_y + 2\frac{\partial \bar{\phi}}{\partial d_z} \cdot d_x d_z + \bar{\phi} \cdot d_x \right\} \cdot d_r$$

$$E_y = \left\{ 2\frac{\partial \bar{\phi}}{\partial d_x} \cdot d_y d_x - \frac{\partial \bar{\phi}}{\partial d_y} (d_r^2 - 2d_y^2) + 2\frac{\partial \bar{\phi}}{\partial d_z} \cdot d_y d_z + \bar{\phi} \cdot d_y \right\} \cdot d_r$$

$$E_z = \left\{ 2\frac{\partial \bar{\phi}}{\partial d_x} \cdot d_z d_x + 2\frac{\partial \bar{\phi}}{\partial d_y} \cdot d_z d_y - \frac{\partial \bar{\phi}}{\partial d_z} \cdot (d_r^2 - d_z^2) + \bar{\phi} \cdot d_z \right\} \cdot d_r$$

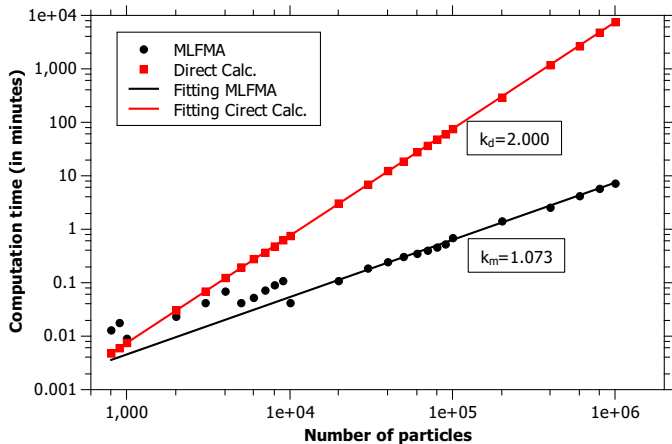
with

$$d_r = \sqrt{d_x^2 + d_y^2 + d_z^2}.$$

## Description of the MLFMA

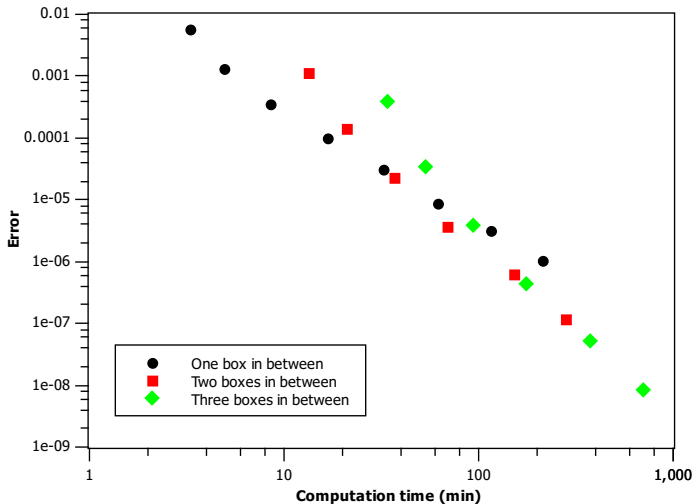
- **Construct** the hierarchical box structure (**partial tree**).
- **Upwards**: Calculate the multipole expansions for all the boxes.
- **Downwards**: For each box, check its the relation with other boxes and operate according to the above table. Then translate the local expansion from parent boxes to the child boxes.
- **Calculate the potential/field**, which comes from direct calculation and multipole or local expansions.

## Compare the MLFMA with direct calculation

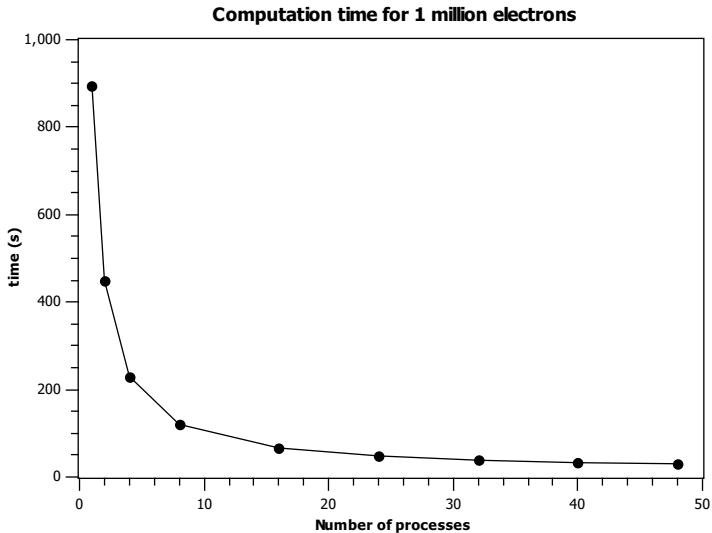




## Accuracy and computation time

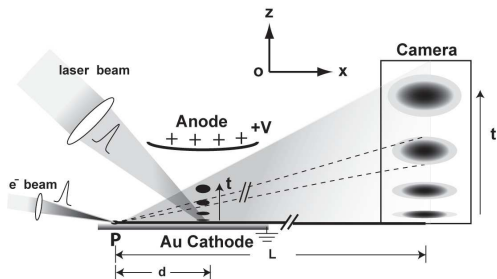


# Numerical experiments



## II. Simulation results on femtosecond imaging and spectroscopy

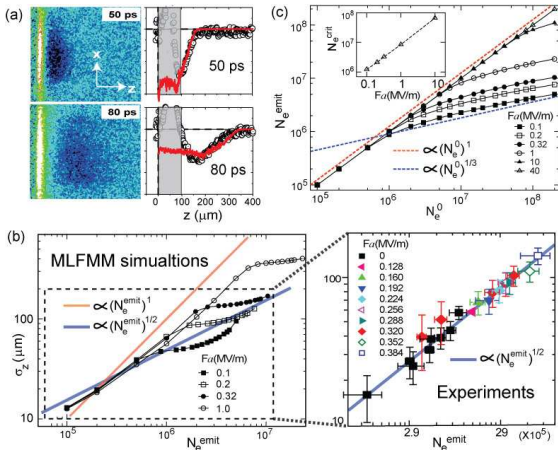
Experiment setup:



Simulation model:

- Three-step model of photoemission
- Extraction field  $F_a$ , space charge field, surface image charges
- Up to several millions of macro-particles in simulation

# Simulation

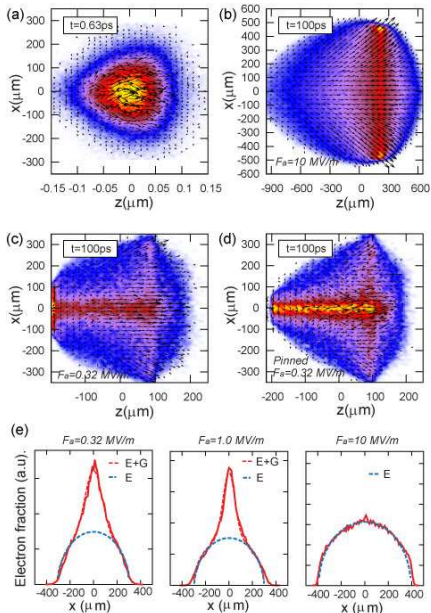


(a) Electron bunch profiles (b) Bunch length vs.  $N_e^{\text{emit}}$   
 (c)  $N_e^{\text{emit}}$  vs.  $N_e^0$  showing virtual cathode effect

1 .Z. Tao, H. Zhang, P. M. Duxbury, M. Berz, C.-Y. Ruan, Journal of Applied Physics 111 (2012) 044316

2. J. Portman, H. Zhang, Z. Tao, etc., Applied Physics Letter, in press, 2013

# Simulation



(a) Initial electron bunch profiles in x-z plane

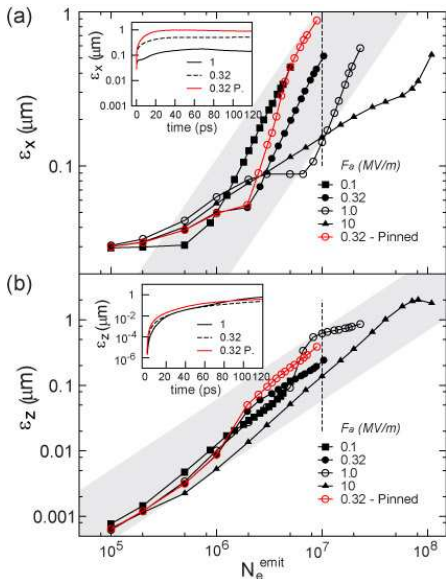
(b) Electron bunch profiles in x-z plane at  $100\text{ps}$  under strong  $F_a$  ( $10\text{ MV/m}$ )

(c) Electron bunch profiles in x-z plane at  $100\text{ps}$  under weak  $F_a$  ( $0.32\text{ MV/m}$ )

(d) Electron bunch profiles in x-z plane at  $100\text{ps}$  under weak  $F_a$  ( $0.32\text{ MV/m}$ ) with a pinned image field

(e) Electron charge distributions under different fields

# Simulation

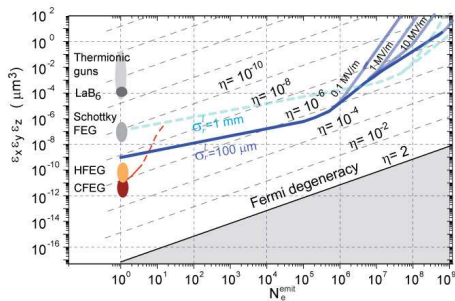


(a)  $\varepsilon_x$  for different field and different  $N_e^{\text{emit}}$ . Affected by the virtual cathode effect.

(b)  $\varepsilon_y$  for different field and different  $N_e^{\text{emit}}$ . Not affected by the virtual cathode effect. Driven by internal field and nonlinearities.

# Simulation

Figure: 6D emittance  $\varepsilon_x \varepsilon_y \varepsilon_z$  versus  $N_e^{\text{emit}}$  for the extended electron sources with sizes  $\sigma_r^l$  (100  $\mu\text{m}$ , 1mm), thermionic guns, Schottky, cold (CFEG) and heated (HFEG) field-emission guns<sup>1</sup>



<sup>1</sup> Science of Microscopy, Eds. P.W. Hawkes, and J.C.H. Spence (Springer, New York, 2008)

- Degeneracy ( $\eta = B_{6D}\varepsilon_0$ ,  $B_{6D} = N_e/(\varepsilon_x^2\varepsilon_z)$ ) can be improved with flat photoemission cathode until the virtual cathode threshold.
- Increasing extraction field is helpful, because it defers the onset of virtual cathode effect.
- Large emitting area increases the  $N_e^{\text{emit}}$ , but not necessarily degeneracy (brightness).
- No gain by using sharp emitters (FEGs or atom-sized emitters) due to their poor emittance scaling with  $N_e$ .



## 1. The multiple Level Fast Multipole Algorithm

- Grid-free, works for any arbitrary charge distribution with an efficiency  $O(N)$ .
- Calculate both the potential/field and its derivatives

## 2. Studied the virtual cathode effect quantitatively in femtosecond electron generation

- $N_e^{\text{emit}}$ ,  $\varepsilon_x$ ,  $\eta$
- Increase  $F_a$  can defer the onset of virtual cathode effect

## 3. Transition into virtual cathode regime might be described by two fluids

## 4. $\eta$ is affected by different factors.



**THANK  
YOU!**